

Landau damping and Plasma echoes

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Of interest in Plasma Physics:

- The dynamics of charged particles (ions / electrons)
- The classical Vlasov-Poisson system¹ on $\mathbb{T}^d \times \mathbb{R}^d$, $d \geq 1$:

$$\partial_t f + v \cdot \nabla_x f + E \cdot \nabla_v f = 0, \quad \nabla_x \cdot E = \rho - 1$$

for density distribution $f(t, x, v) \geq 0$ and charge density $\rho = \int f dv$.

- The neglecting of
 - Boltzmann/Landau collision (short range) \implies Vlasov
 - Magnetic effect (Maxwell at $c \rightarrow \infty$) \implies Poisson
 - The dynamics of ions (mass ratio $m_e/m_i \rightarrow 0$) \implies Electrons

¹that is, transport along $\dot{x} = v$, $\dot{v} = E(t, x)$

(in great similarity to 2D Euler....)

- Hamiltonian:

$$\mathcal{H}[f] = \frac{1}{2} \iint |v|^2 f \, dx dv + \frac{1}{2} \int |E|^2 \, dx$$

- Invariant Casimir's:

$$\mathcal{C}[f] = \iint \Phi(f) \, dx dv.$$

(e.g., $\int \rho \, dx = 1$ for all times; perfect for $\nabla_x \cdot E = \rho - 1$ on \mathbb{T}^d)

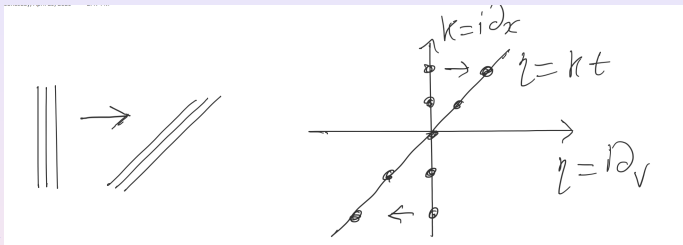
- The global Cauchy problem is classical: e.g., Glassey (SIAM, '96)

- **Large time behavior** of spatially homogenous equilibria

$$f = \mu(v), \quad E = 0, \quad \rho = 1$$

- In the large times,
 - **Landau damping**: decay of the electric field
 - **Plasma echoes**: new waves are excited at later times...
- Linear damping: Landau '46.
- Progress on Nonlinear damping: Mouhot-Villani ('11), Bedrossian-Masmoudi-Mouhot ('16), Bedrossian ('16), Lin Zeng ('11).

- Phase mixing: $\partial_t f + v \partial_x f = 0$.



Explicitly,

$$f_0(x, v) = g_0(v) e^{ikx + i\eta v} \implies f(t, x, v) = g_0(v) e^{ikx + i(\eta - kt)v}$$

$$\rho_0(x) = \hat{g}_0(\eta) e^{ikx} \implies \rho(t, x) = \hat{g}_0(\eta - kt) e^{ikx}$$

Regularity in $v \implies$ Decay of the electric field $E = \frac{1}{ik} \rho$.

(analyticity) \implies (exponentially localized near $t = \frac{\eta}{k}$)

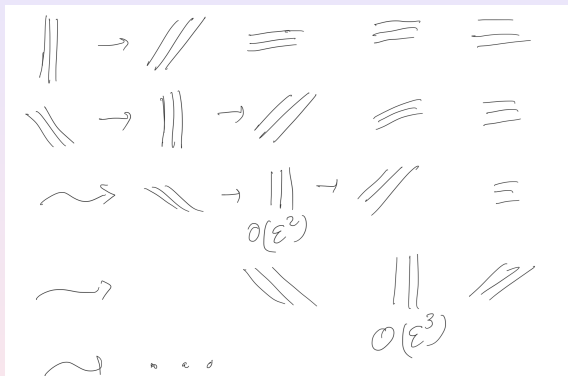
- Plasma echoes: $\partial_t f + v \partial_x f + E \partial_v f = 0, \quad \partial_x E = \rho - 1$

$\hat{E}_{k_1}(t) \lesssim \varepsilon e^{-\langle \gamma_1 - k_1 t \rangle}$
 $\hat{E}_{k_2}(t) \lesssim \varepsilon e^{-\langle \gamma_2 - k_2 t \rangle}$
 $O(\varepsilon^2)$
 $t \sim \frac{\gamma_1 + \gamma_2}{k_1 + k_2}$
 $\gg t_{k_1, \gamma_1}, t_{k_2, \gamma_2}$

→ Interaction through $E^1 \partial_v f^2$
Echo: → 3rd wave
 → 4th wave
 → ...

what's going on with 1?

A cascade of echoes?



Actually, data is **large in any Sobolev**:

$$f^0(x, v) = \mu(v) + f_1^0(v)e^{iK_1x+iL_1v} + f_2^0(v)e^{iK_2x+iL_2v}$$

Observations:

- Electric field is localized near critical times (exponentially, if analytic).
- Echoes “take time” to appear, and so not seen in analytic (or Gevrey) classes studied by Mouhot-Villani.

- Linear Landau damping:

$$\partial_t f + v \partial_x f + E \partial_v \mu = 0, \quad \partial_x E = \rho,$$

yielding a *closed* equation for density:

$$\hat{\rho}_k(t) + \int_0^t \hat{E}_k(s) \widehat{\partial_v \mu}(k(t-s)) ds = \hat{f}_{k,kt}^0$$

$$\hat{\rho}_k(t) + [t\hat{\mu}(kt)] *_t \hat{\rho}_k(t) = \hat{f}_{k,kt}^0 \implies \mathcal{L}[\hat{\rho}_k(t)] = \frac{\mathcal{L}[\hat{f}_{k,kt}^0]}{1 + \mathcal{L}[t\hat{\mu}(kt)]}$$

assuming Penrose. Further, writing the resolvent kernel:

$$\frac{1}{1 + \mathcal{L}[t\hat{\mu}(kt)]} = 1 - \frac{\mathcal{L}[t\hat{\mu}(kt)]}{1 + \mathcal{L}[t\hat{\mu}(kt)]} \quad (1 = \text{phase mixing})$$

Theorem (Linear Landau damping: Grenier-Toan-Rodnianski)

Assume Penrose: $|1 + \mathcal{L}[t\hat{\mu}(kt)](\lambda)| \gtrsim 1$. Then,

$$\hat{\rho}_k(t) = \hat{S}_k(t) + \int_0^t \hat{K}_k(t-s)\hat{S}_k(s) ds, \quad |\hat{K}_k(t)| \lesssim e^{-\theta_0|kt|},$$

with $\hat{S}_k(t) = \hat{f}_{k,kt}^0$. Namely, *Landau \approx phase mixing* (under Penrose).

- Linear damping:

$$\hat{\rho}_k(t) \lesssim \hat{f}_{k,kt}^0 + \dots \lesssim \begin{cases} \langle kt \rangle^{-\sigma}, & \text{Sobolev data} \\ e^{-\langle kt \rangle^\gamma}, & \text{Analytic or Gevrey-}\gamma \text{ data} \end{cases}$$

- Penrose holds for monotone equilibria (or even small bumps in tail), including Gaussians and any radial positive 3D or higher-D equilibria.

- Generator functions:

$$G[g](z) := \sum_{k \in \mathbb{Z}} \int_{\mathbb{R}} e^{z \langle k, \eta \rangle} \left[|\widehat{g}_{k, \eta}| + |\partial_{\eta} \widehat{g}_{k, \eta}| \right] \langle k, \eta \rangle^{\sigma} d\eta$$

for analyticity radius $z \geq 0$ (eventually $z = \lambda(t)$):

- $G[\partial_x g] \leq \partial_z G[g]$ and $G[\partial_v g] \leq \partial_z G[g]$
- Derving simple transport inequality:

$$\partial_t G[g] \leq A[g] + B[g] \partial_z G[g].$$

- As for Gevrey- γ : use L^2 with weight $e^{z \langle k, \eta \rangle^{\gamma}}$ and integrate by parts.
- **Application:** Instability of boundary layers, Inviscid limit,...

- **Transport equation** for $g(t, x, v) = f(t, x + vt, v)$:

$$\partial_t g = -E(t, x + vt) \partial_v \mu(v) - E(t, x + vt) (\partial_v - t \partial_x) g.$$

- **Transport inequality** for norm $G[g](t, z)$:

$$\partial_t G[g] \lesssim F[E] + (1 + t) F[E] \partial_z G[g]$$

where

$$F[E](t, z) := \sum_{k \in \mathbb{Z}} e^{z \langle k, kt \rangle} |\hat{E}_k(t)| \langle k, kt \rangle^\sigma.$$

$$\partial_t G[g] \lesssim F[E] + (1+t)F[E]\partial_z G[g]$$

- **Extra fast decay** on $F[E]$ controls the shrinking of analyticity radius:

$$\lambda'(t) + C_0(1+t)F[E](t, \lambda(t)) \leq 0 \implies \partial_t G[g](t, \lambda(t)) \lesssim F[E](t, \lambda(t)).$$

Theorem (Nonlinear Landau damping: Mouhot-Villani)

For small Gevrey-3⁻ initial data near Penrose stable equilibria:

$$E(t, x) \rightarrow 0, \quad f(t, x + vt, v) \rightarrow f_\infty(x, v),$$

exponentially fast in $\langle t \rangle^\gamma$.

- Again, a “closed” equation for density,

$$\widehat{E}_k(t) + [t\widehat{\mu}(kt)] \star_t \widehat{E}_k(t) = \widehat{S}_k(t)$$

$$\begin{aligned} \widehat{S}_k(t) &:= \frac{1}{ik} \widehat{f}_{k,kt}^0 - \sum_{l \neq 0} \int_0^t (t-s) \widehat{g}_{k-l,kt-ls}(s) \widehat{E}_l(s) ds \\ &\lesssim k^{-1} \widehat{f}_{k,kt}^0 + \sum_{l \neq 0} \int_0^t (t-s) \frac{e^{-\lambda(s)\langle k-l, kt-ls \rangle^\gamma}}{\langle kt-ls \rangle^\sigma} \frac{e^{-\lambda(s)\langle l, ls \rangle^\gamma}}{\langle ls \rangle^\sigma} e^{-\langle s \rangle^{\gamma-\delta}} ds \\ &\lesssim e^{-\lambda(t)\langle k, kt \rangle^\gamma} \langle kt \rangle^{-\sigma} e^{-\langle t \rangle^{\gamma-\delta}} \quad (\text{extra decay for } E) \end{aligned}$$

Try analyticity radius $\lambda(t) = \lambda_0 + t^{-\delta}$. For s away from t ,

$$e^{-(\lambda(s)-\lambda(t))\langle k, kt \rangle^\gamma} \leq e^{-\theta_0 |t-s|\langle k, kt \rangle^\gamma / t^{1+\delta}} \leq e^{-\theta'_0 \langle t \rangle^{\gamma-\delta}}.$$

Suppression of echoes: For $|kt - ls| \ll t$,

$$k(t - s) = kt - ls + (l - k)s \gtrsim |l - k|t \gtrsim t$$

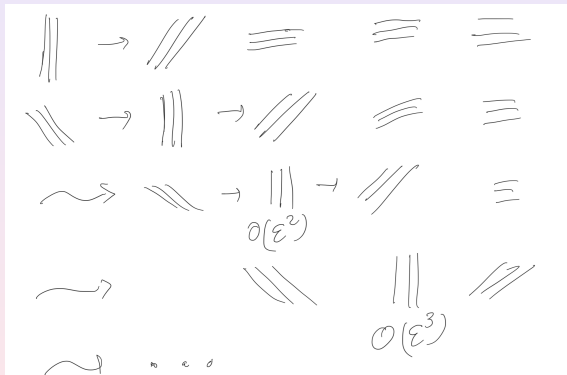
and so

$$e^{-(\lambda(s) - \lambda(t))\langle k, kt \rangle^\gamma} \leq e^{-\theta_0 |k|(t-s)\langle kt \rangle^\gamma / |k|t^{1+\delta}} \leq e^{-\theta'_0 \langle kt \rangle^\gamma / |k|t^\delta}$$

$$\int_0^t (t - s) \langle kt - ls \rangle^{-\sigma} ds \lesssim \frac{t}{k^2} = \frac{|kt|}{k^3} \lesssim \frac{|k|^{\frac{1-\delta}{\gamma-\delta}}}{|k|^3} \lesssim 1$$

provided $3\gamma > 1$, giving Mouhot-Villani's for Gevrey-3⁻.

An infinite cascade of echoes:



- Highly oscillatory data:

$$f^0(x, v) = \mu(v) + \sum_{(k, \eta) \in \mathbb{Z} \setminus \{0\} \times \mathbb{Z}} \epsilon f_{k, \eta}^0(v) e^{iKkx + iL\eta v}$$

for large K, L and small ϵ (hence, arbitrarily large in any Sobolev).

- Search for a ***complete*** echo cascade Ansatz:

$$f(t, x, v) = \mu(v) + \sum_{p=1}^{\infty} \epsilon^p \sum_{(k, \eta) \in \mathbb{Z} \times \mathbb{Z}} f_{k, \eta, p}(t, v) e^{iKkx + i(L\eta - Kkt)v}$$

where $f_{k, \eta, p}(t, v)$ solves the linearized Vlasov-Poisson, inductively in $p \geq 1$.

- **Application:** a complete instability of boundary layers.

- Assumptions:

- Echoes times are uniformly bounded:

$$L \lesssim K.$$

- Coefficients are analytic:

$$|\widehat{f}_{k,\eta}^0(\eta')| \leq e^{-2\lambda_0 \langle k,\eta,\eta' \rangle}, \quad \forall k, \eta, \eta'.$$

Proposition (Estimates for echoes)

There is some universal constant C_0 so that

$$\begin{aligned} |\widehat{f}_{k,\eta,p}(t, \eta')| &\leq C_0^p e^{-\lambda_p(t)\langle k,\eta,p,\eta' \rangle}, \\ |\widehat{E}_{k,\eta,p}(t)| &\leq C_0^p e^{-\lambda_p(t)\langle k,\eta,p,L\eta-Kkt \rangle} \langle t \rangle^{-\sigma}, \end{aligned}$$

where “analyticity” radius $\lambda_p(t)$ is defined by

$$\lambda_p(t) = \lambda_0 + \langle t \rangle^{-\delta} + p^{-\delta}, \quad \delta \ll 1.$$

Proof.

It follows from Landau damping and “localized” interaction. □

Theorem (An infinite cascade of echoes: Grenier-Toan-Rodnianski)

There exists a complete echo solution to the Vlasov-Poisson system:

- initial data are arbitrarily large in $W^{s,\infty}$ Sobolev, for any $s > 0$
- Landau damping holds for echoes:

$$E(t, x) = \sum_{(k, \eta, p) \in \mathbb{Z} \times \mathbb{Z} \times \mathbb{N}^*} \epsilon^p \widehat{f}_{k, \eta, p}(t, Kkt - L\eta) \frac{e^{iKkx}}{iKk} \longrightarrow 0$$

exponentially fast in any Sobolev spaces $W^{s,q}$, $s \geq 0, q \geq 1$.